

On the implementation of X-LES in a high-order implicit DG solver

F. Bassi¹, L. Botti¹, A. Colombo¹, A. Crivellini², A. Ghidoni³,
M. Lorini³, F. Massa¹, G. Noventa³

¹Università degli studi di Bergamo

²Università politecnica delle Marche

³Università degli studi di Brescia

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...with the contribution of

Francesco Bassi¹

Alessandro Colombo¹

Lorenzo Botti¹

Francesco Carlo Massa¹

Marco Savini¹

Nicoletta Franchina¹

Antonio Ghidoni²

Gianmaria Noventa²

Marco Lorini²

Stefano Rebay²

Andrea Crivellini³

Carmine De Bartolo⁴

Alessandra Nigro⁴

Daniele Di Pietro⁵

Pietro Tesini⁶

¹ Università degli Studi di Bergamo

² Università degli Studi di Brescia

³ Università Politecnica delle Marche

⁴ Università della Calabria

⁵ University of Montpellier, France

⁶ SKF, Sweden

The logo for TILDA features the word "TILDA" in a bold, black, sans-serif font. The letter "I" is replaced by a red square. A thick black arrow points to the right, starting from the top of the "I" and extending above the "A".

**Towards Industrial LES/DNS in Aeronautics
Paving the Way for Future Accurate CFD
grant agreement No.635962**

Brief code summary

- Discontinuous Galerkin (DG) method on hybrid grids
- Physical frame orthonormal basis functions
- 2D/3D steady and unsteady compressible and incompressible flows
- Explicit and implicit time accurate integration
- Fixed or rotating frame of reference
- Euler
- Navier–Stokes
- RANS coupled with the k - ω (EARSM)
- Hybrid RANS/LES (X-LES)
- MPI parallelism
- Fortran language

Brief code summary

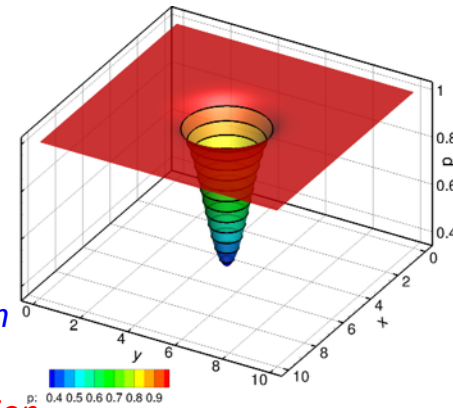
- Discontinuous Galerkin (DG) method on hybrid grids
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- Explicit and **implicit time accurate integration - variables**
- Fixed or rotating frame of reference
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- **Hybrid RANS/LES (X-LES)**
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Implicit accurate time integration

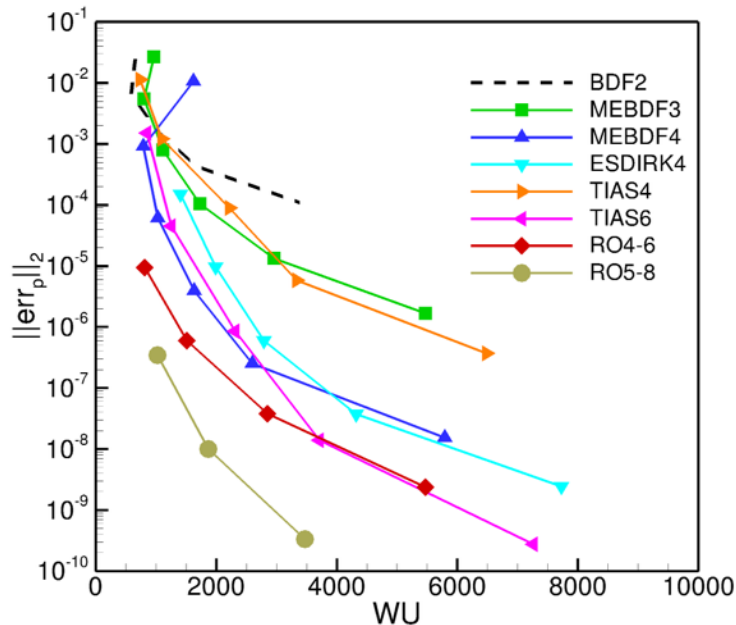
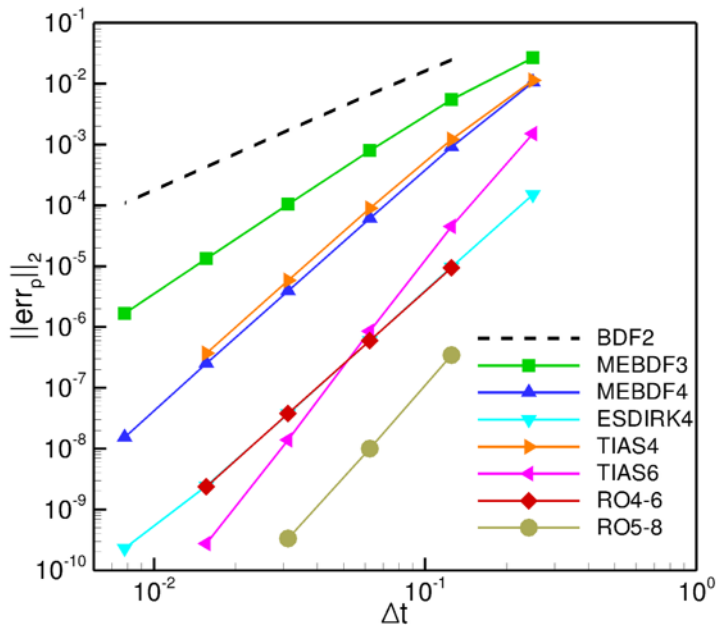
Several high-order temporal schemes are implemented

- Modified Extended BDF
 - Two Implicit Advanced Step-point (TIAS)
 - Explicit Singly Diagonally Implicit R-K (ESDIRK)
 - linearly implicit Rosenbrock method
- }

non-linear systems solution
linear systems solution (here via GMRES)



- i) Hi-O schemes are **more efficient than Lo-O** ones for high required accuracy
- ii) **Rosenbrock-type** schemes are **appealing** both for accuracy and efficiency



Convection of an isentropic vortex P^6 solution on 50X50 el.

Rosenbrock schemes in a nutshell (I/II)

From the DG spatial discretization we obtain a system of non-linear ODEs or DAEs

$$\mathbf{M}_P(\mathbf{W}) \frac{d\mathbf{W}}{dt} + \mathbf{R}(\mathbf{W}) = \mathbf{0} \quad \tilde{\mathbf{R}} = \mathbf{M}_P^{-1} \mathbf{R}$$

$$\mathbf{W}^{n+1} = \mathbf{W}^n + \sum_{j=1}^s m_j \mathbf{Y}_j$$

$$\left(\frac{\mathbf{M}_P}{\gamma \Delta t} + \mathbf{J} - \frac{\partial \mathbf{M}_P}{\partial \mathbf{W}} \tilde{\mathbf{R}} \right)^n \mathbf{Y}_i = -\mathbf{M}_P^n \left[\tilde{\mathbf{R}} \left(\mathbf{W}^n + \sum_{j=1}^{i-1} a_{ij} \mathbf{Y}_j \right) - \sum_{j=1}^{i-1} \frac{c_{ij}}{\Delta t} \mathbf{Y}_j \right]$$

$i = 1, \dots, s$

only a **linear system** need to be solved for each stage

i.e. the **Jacobian** $\mathbf{J} = \partial \mathbf{R} / \partial \mathbf{W}$ is **assembled and factored only once per time step**

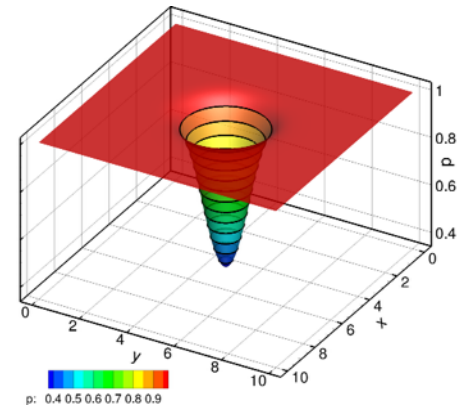
With **orthonormal basis** functions (*physical space*) \mathbf{M}_P reduces to the **identity** for compressible flows with **conservative** variables

For **other sets of variables** their DOFs can be coupled within \mathbf{M}_P thus resulting in a matrix which can **not be diagonal**

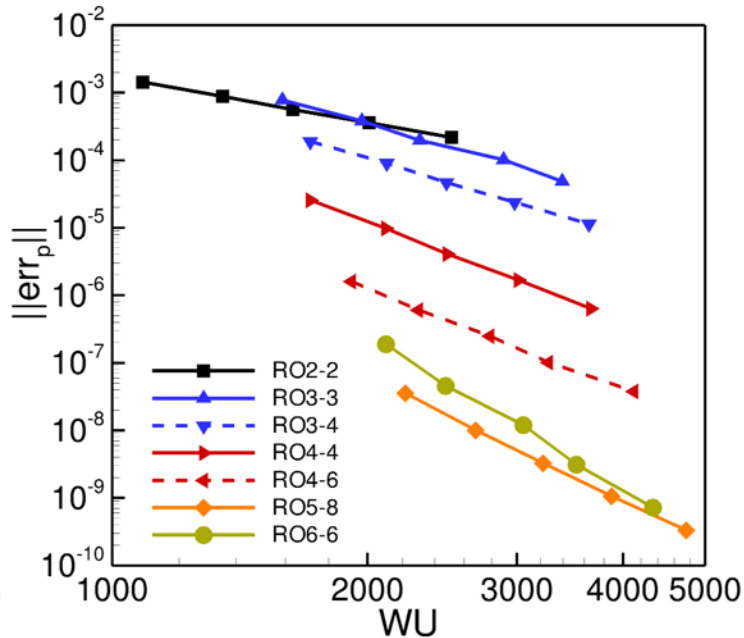
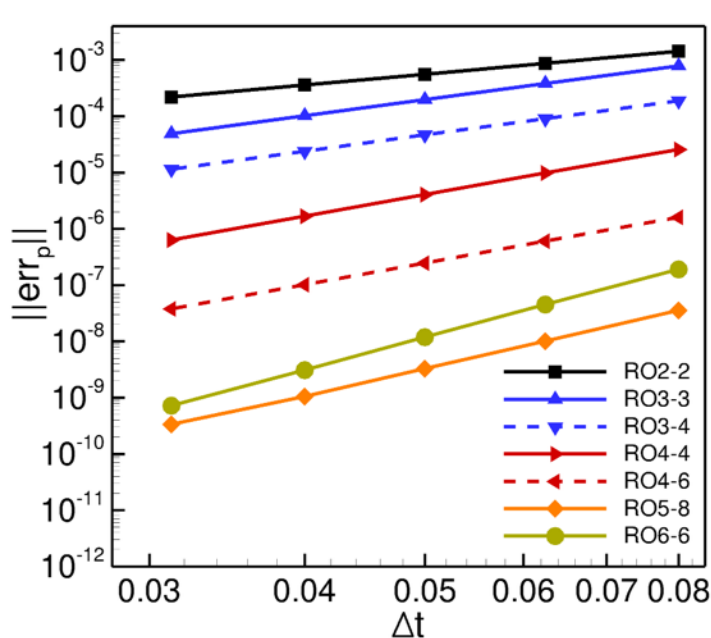
Rosenbrock schemes in a nutshell (II/II)

Several Rosenbrock schemes, from order two to order six, have been compared

No need to “exactly” solve systems: **GMRES tolerance can be increased** with confidence with a significant **reduction of WU**



For a given order of accuracy, among the schemes considered, those with **more stages are more accurate and efficient**, e.g. RO5-8 vs. RO6-6



Convection of an isentropic vortex
 P^6 solution on 50X50 el.

Working variables

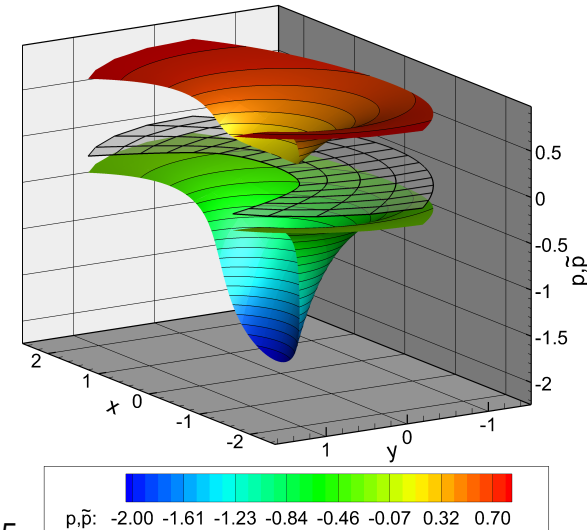
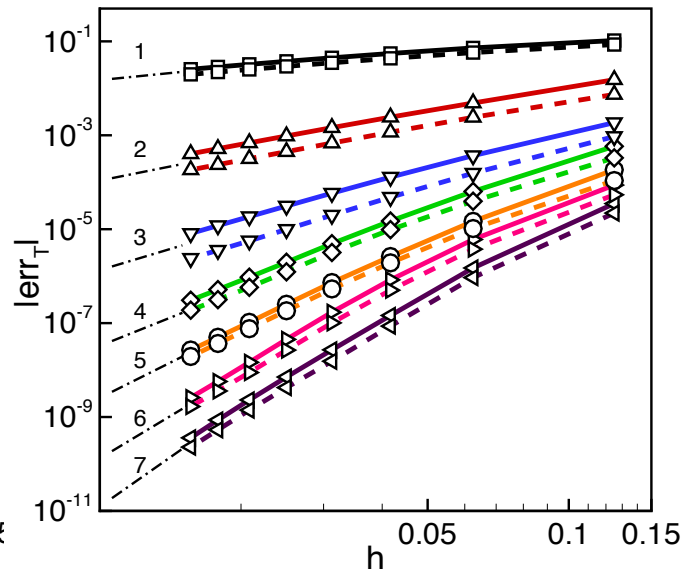
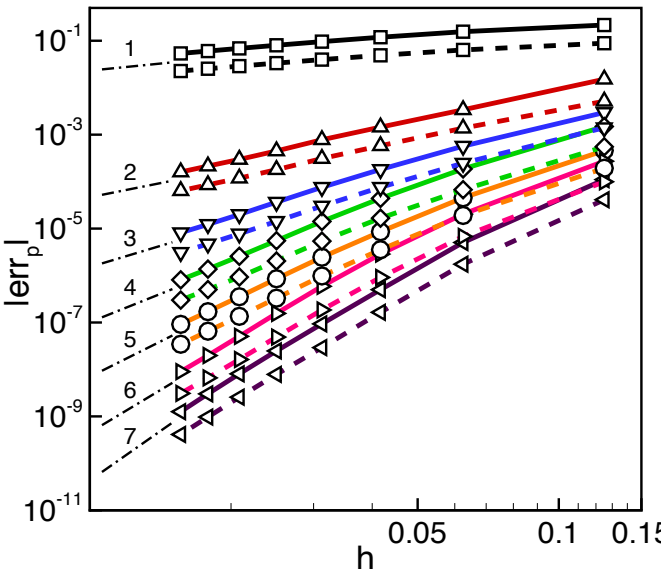
Primitive variables to ensure the positivity of all thermodynamic variables...

We work with **polynomial approximations** not directly for p and T but for **their logarithms** $\tilde{p} = \log(p)$ and $\tilde{T} = \log(T)$

$$\tilde{p}(\mathbf{x}) = \phi_j(\mathbf{x})W_j^{\tilde{p}}, \quad \tilde{T}(\mathbf{x}) = \phi_j(\mathbf{x})W_j^{\tilde{T}}, \quad j = 1, \dots, N_{dof}^k$$

In this way the computed values $p = e^{\tilde{p}}$ and $T = e^{\tilde{T}}$ are always **positive**

Improved code robustness with a low implementation effort (M_P)



eXtra-Large Eddy Simulation (X-LES)

in a nutshell (I/II)

Pros

- hybrid RANS\LES formulation independent from the wall distance
- use in LES mode of a clearly defined SGS based on the k-equation
- use of a k- ω turbulence model integrated to the wall

Cons

the *filter width* parameter is often related to the local element size

$$\frac{\partial}{\partial t}(\rho k) + \frac{\partial}{\partial x_j}(\rho u_j k) = \frac{\partial}{\partial x_j} \left[(\mu + \sigma^* \bar{\mu}_t) \frac{\partial k}{\partial x_j} \right] + P_k - D_k$$

$$\begin{aligned} \frac{\partial}{\partial t}(\rho \tilde{\omega}) + \frac{\partial}{\partial x_j}(\rho u_j \tilde{\omega}) &= \frac{\partial}{\partial x_j} \left[(\mu + \sigma \bar{\mu}_t) \frac{\partial \tilde{\omega}}{\partial x_j} \right] + (\mu + \sigma \bar{\mu}_t) \frac{\partial \tilde{\omega}}{\partial x_k} \frac{\partial \tilde{\omega}}{\partial x_k} \\ &+ P_\omega - D_\omega + C_D \end{aligned}$$

...an “original” interpretation for the X-LES implementation...

Bassi et al. “Time Integration in the Discontinuous Galerkin Code MIGALE - Unsteady Problems”
In IDIHOM: Industrialization of High-Order Methods - A Top-Down Approach, Vol. 128 of
Notes on Numerical Fluid Me- chanics and Multidisciplinary Design Springer International Publishing

eXtra-Large Eddy Simulation (X-LES)

in a nutshell (II/II)

$$\frac{\partial}{\partial t}(\rho k) + \frac{\partial}{\partial x_j}(\rho u_j k) = \frac{\partial}{\partial x_j} \left[(\mu + \sigma^* \bar{\mu}_t) \frac{\partial k}{\partial x_j} \right] + P_k - D_k$$

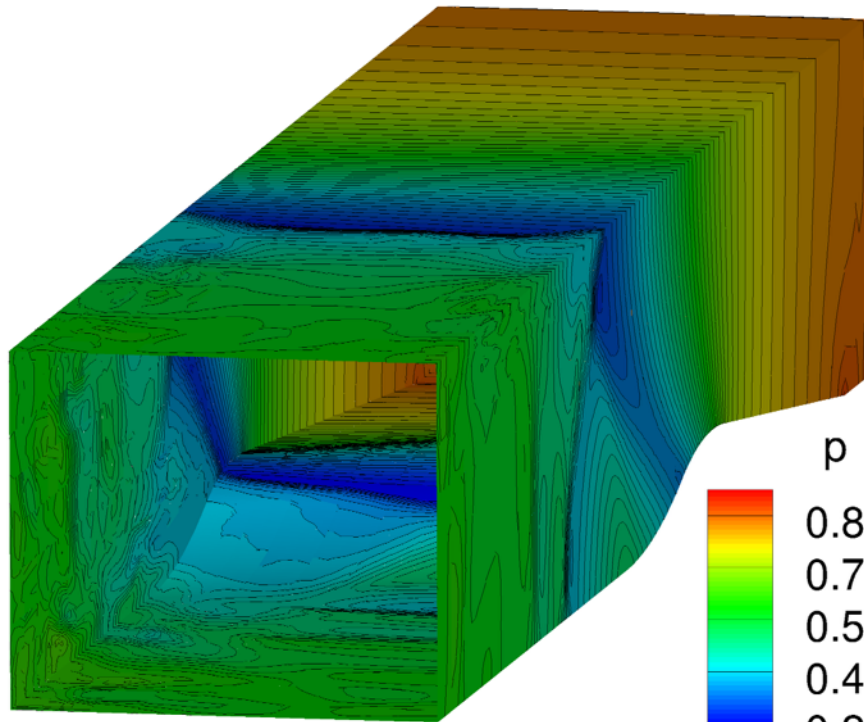
$$\bar{\mu}_t = \alpha^* \frac{\rho \bar{k}}{\hat{\omega}} \quad D_k = \beta^* \rho \bar{k} \hat{\omega} \quad \bar{k} = \max(0, k)$$

$$\hat{\omega} = \max \left(e^{\tilde{\omega}_r}, \frac{\sqrt{\bar{k}}}{C_1 \Delta} \right)$$

	RANS	LES	ILES
$\bar{\mu}_t$	$\alpha^* \frac{\rho \bar{k}}{e^{\tilde{\omega}_r}}$	$\alpha^* \rho \sqrt{\bar{k}} C_1 \Delta$	0
D_k	$\beta^* \rho \bar{k} e^{\tilde{\omega}_r}$	$\beta^* \rho \frac{\bar{k}^{\frac{3}{2}}}{C_1 \Delta}$	0

X-LES of a shock BL interaction on a swept bump (AR2 of HiOCFD, TC-F3 of TILDA)

P² converged computations with **RANS+k- ω** (also in its low-Re version) and **EARSM1** have been performed and used as **initialization** for **X-LES**

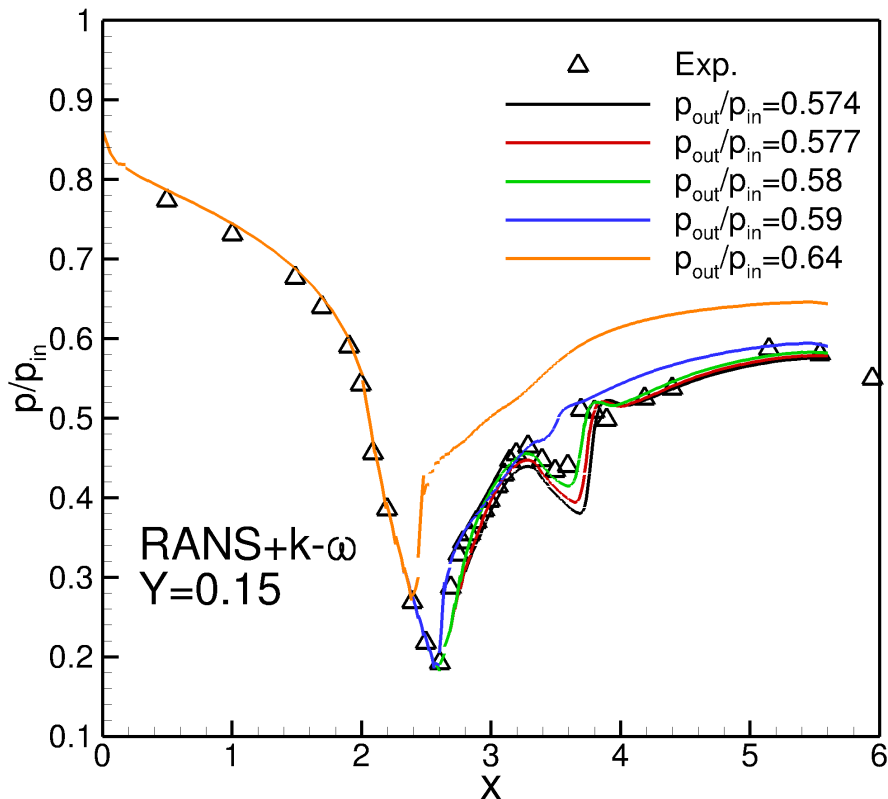


72960 hexahedral elements

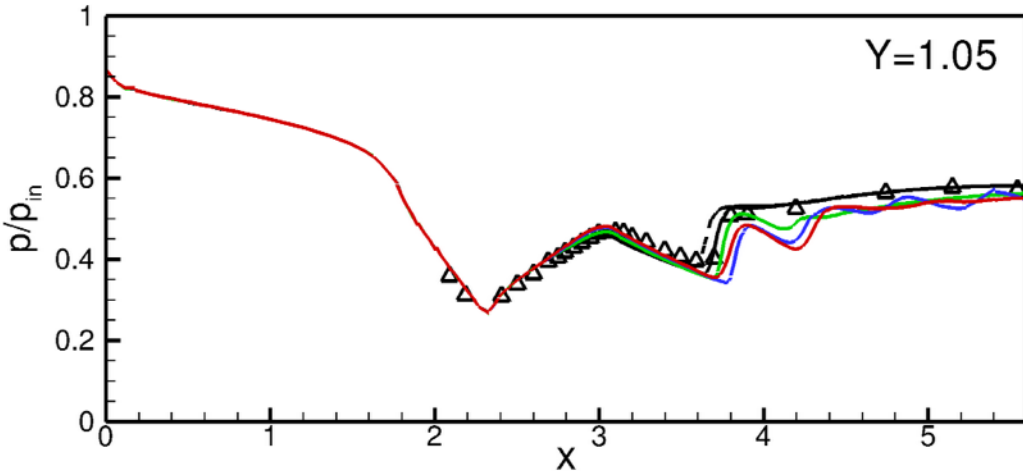
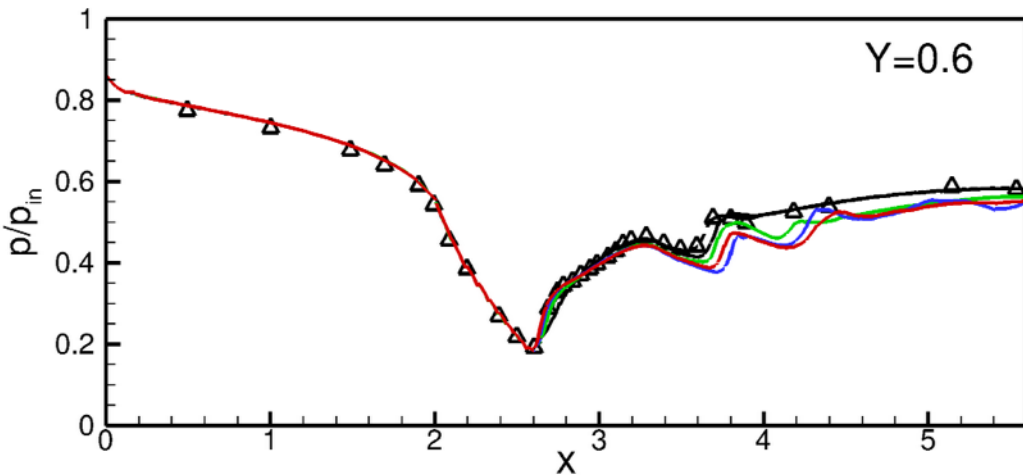
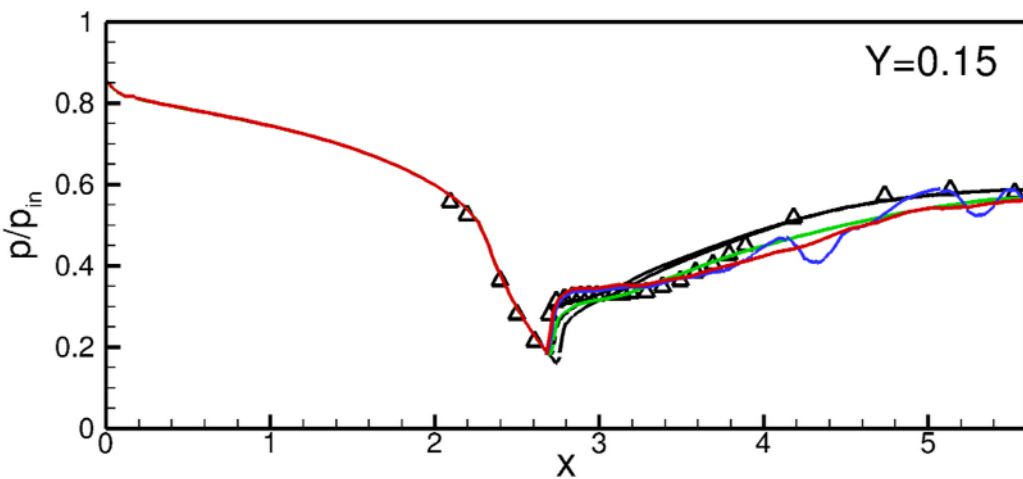
- Inlet boundary conditions
 - $p_{0i} = 92000\text{Pa}$
 - $T_{0i} = 300\text{K}$
 - $Re_H = 1.69 \times 10^6$
- Outlet static pressure used to impose the shock position (*model dependent!*)
- LBE(RO1-1) to quickly find the “right” pressure ratio
- RO3-3 for the time-accurate solution
- Filter width $\Delta = 5e-2$ (***strong influence!***)
- parallel computation on 96 cores

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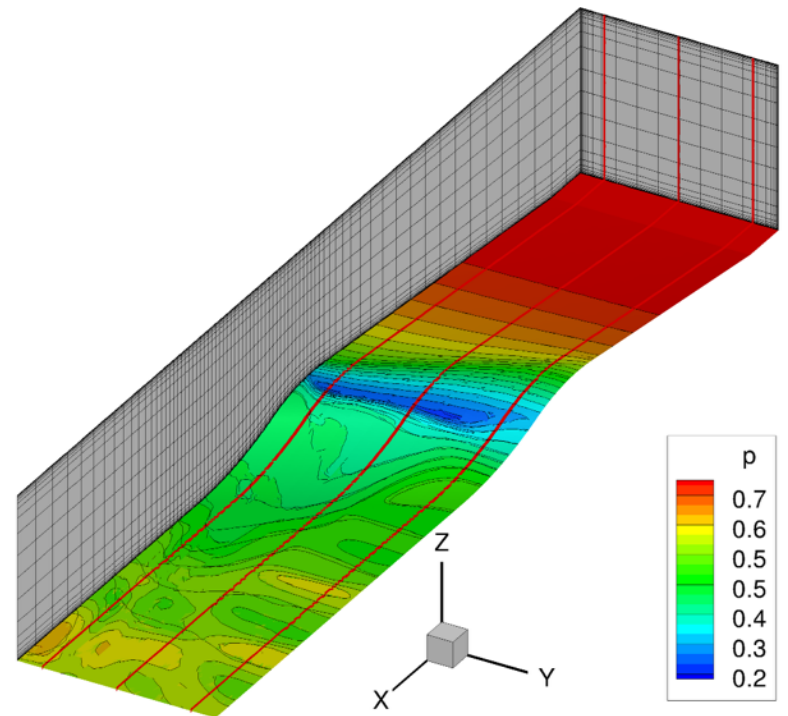
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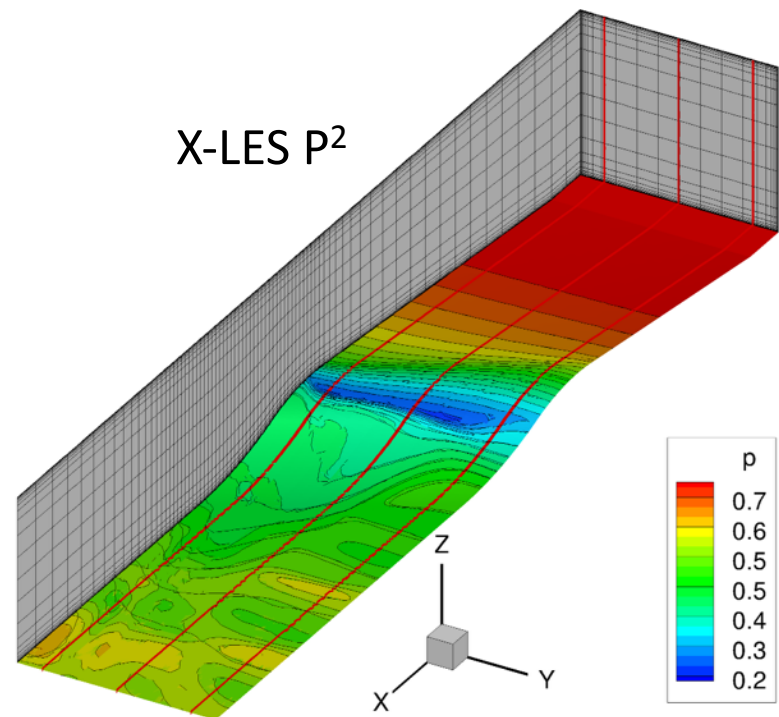
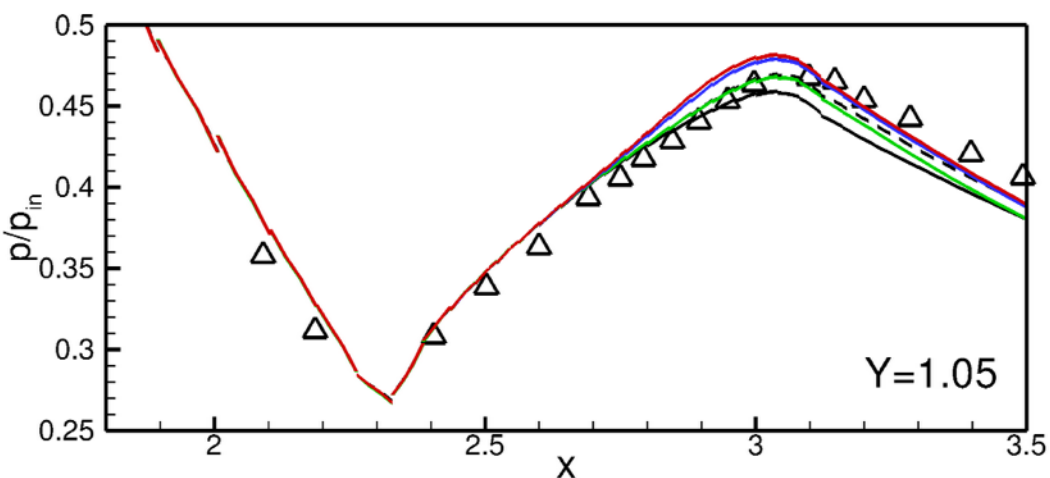
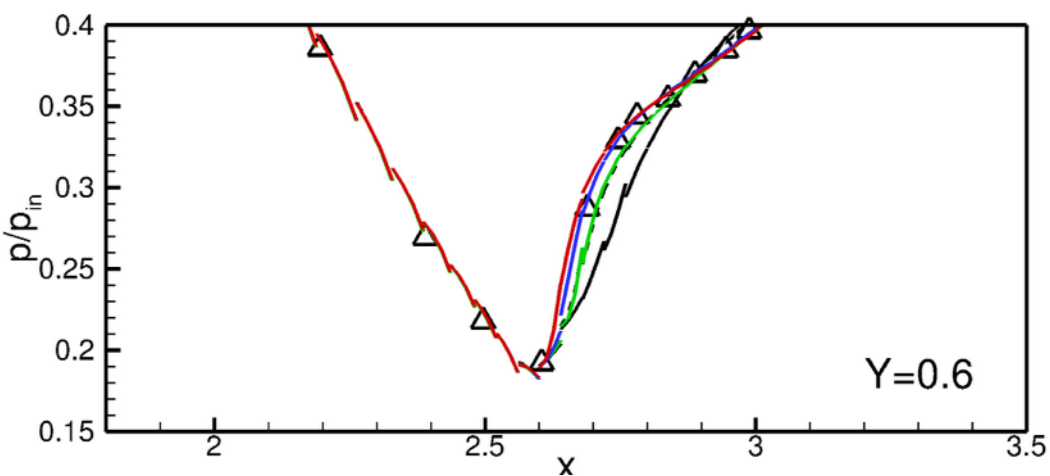
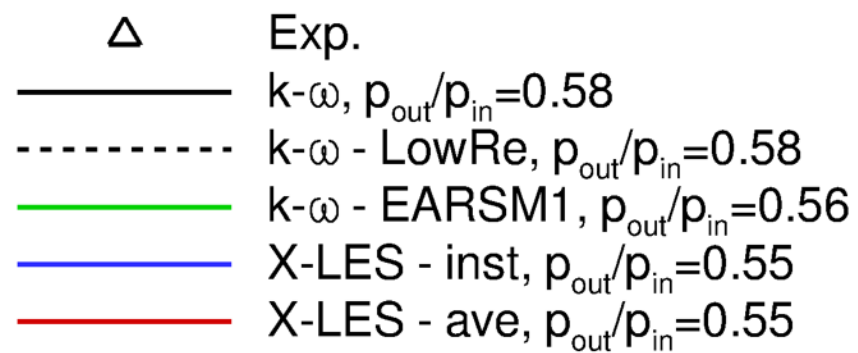
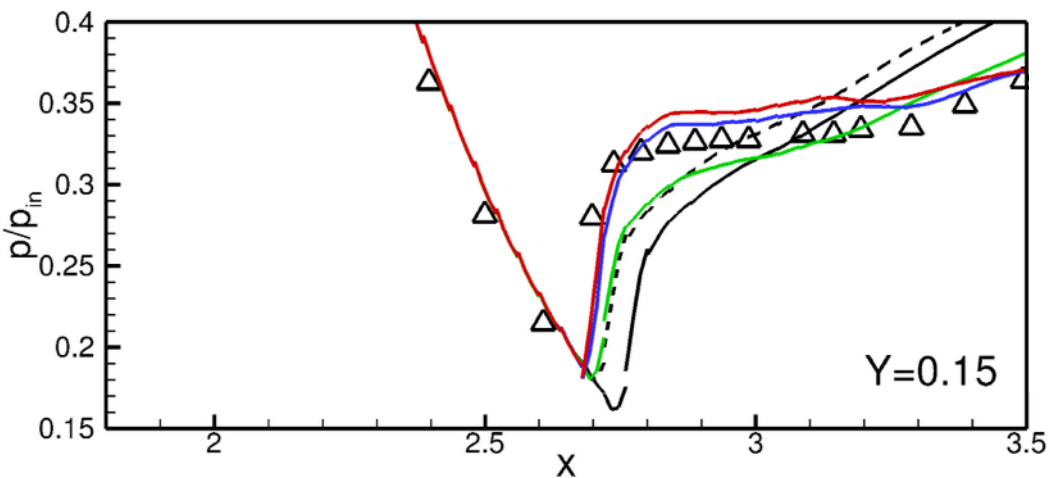
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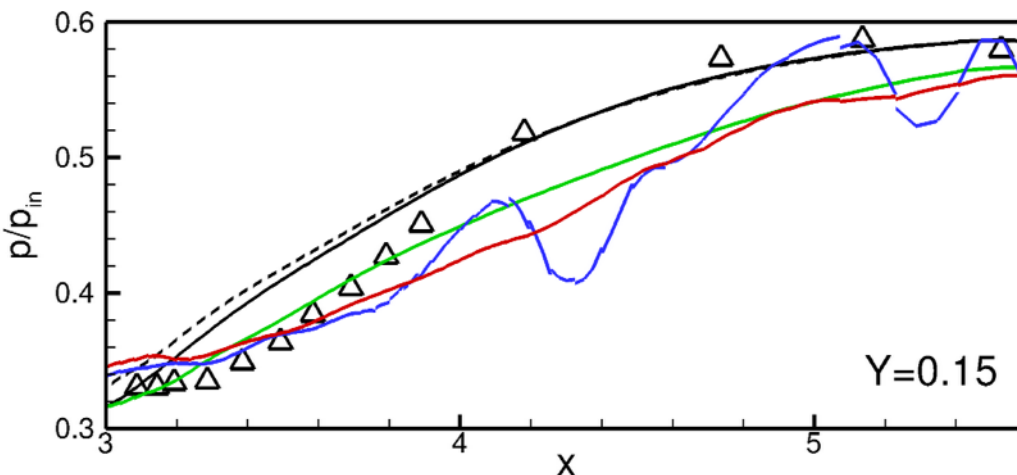
- Δ Exp.
- k- ω , $p_{out}/p_{in}=0.58$
- - - k- ω - LowRe, $p_{out}/p_{in}=0.58$
- k- ω - EARSM1, $p_{out}/p_{in}=0.56$
- X-LES - inst, $p_{out}/p_{in}=0.55$
- X-LES - ave, $p_{out}/p_{in}=0.55$



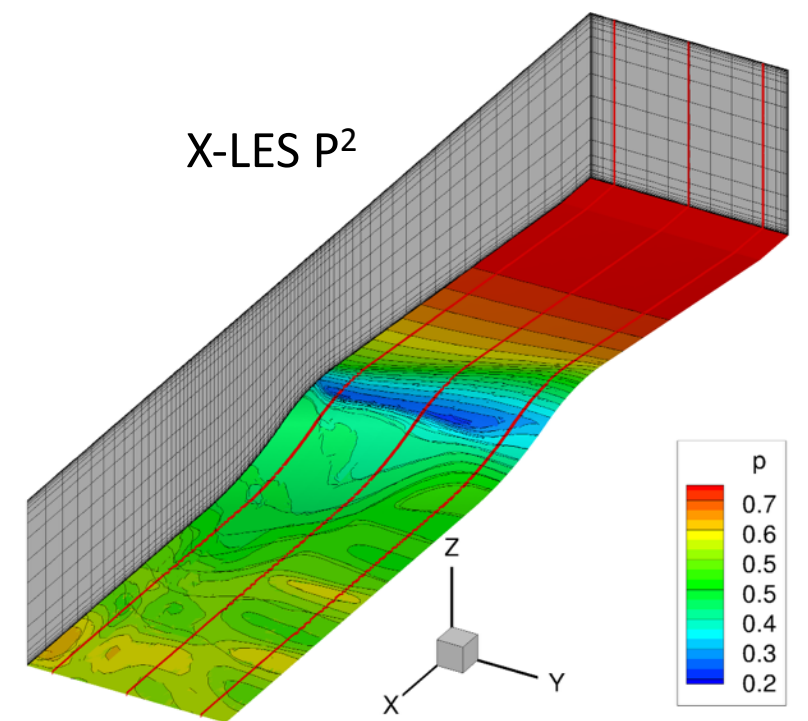
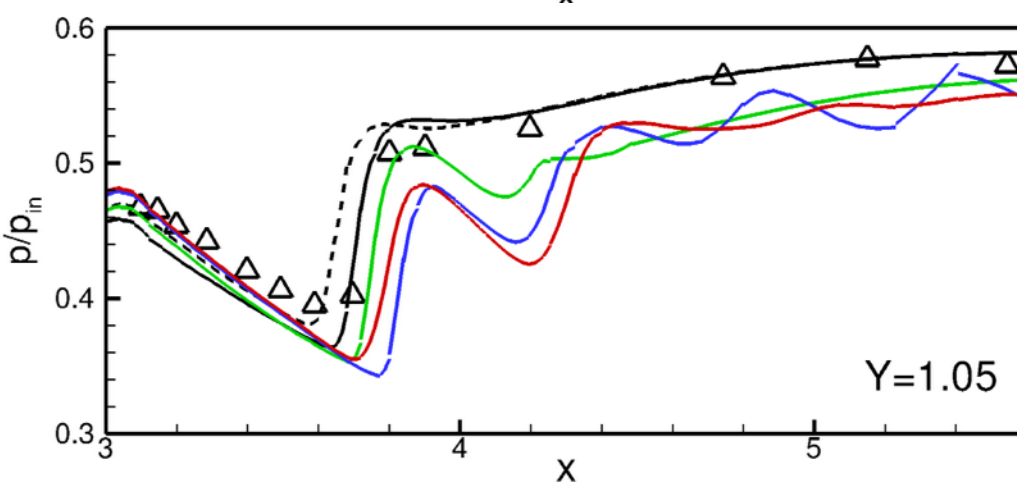
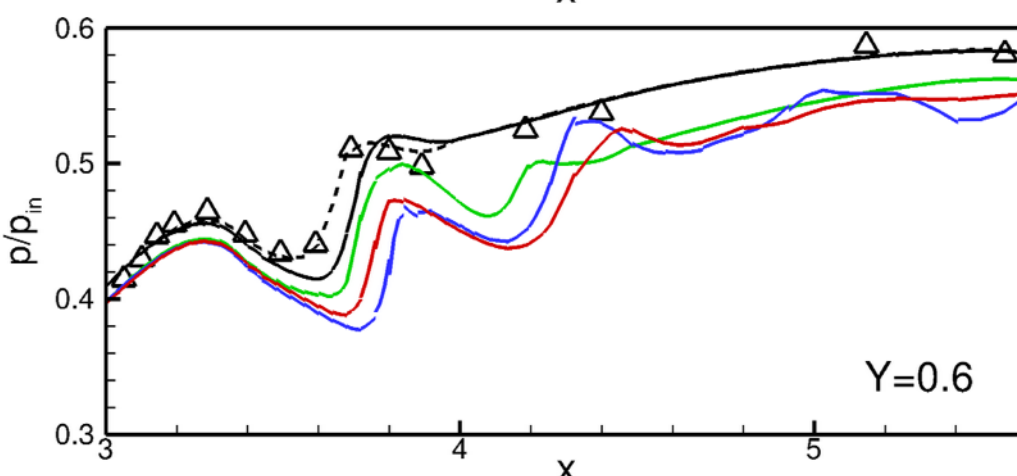
The profiles are in reasonable agreement with the experiments and the numerical results of Cahen et al.



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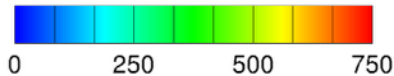
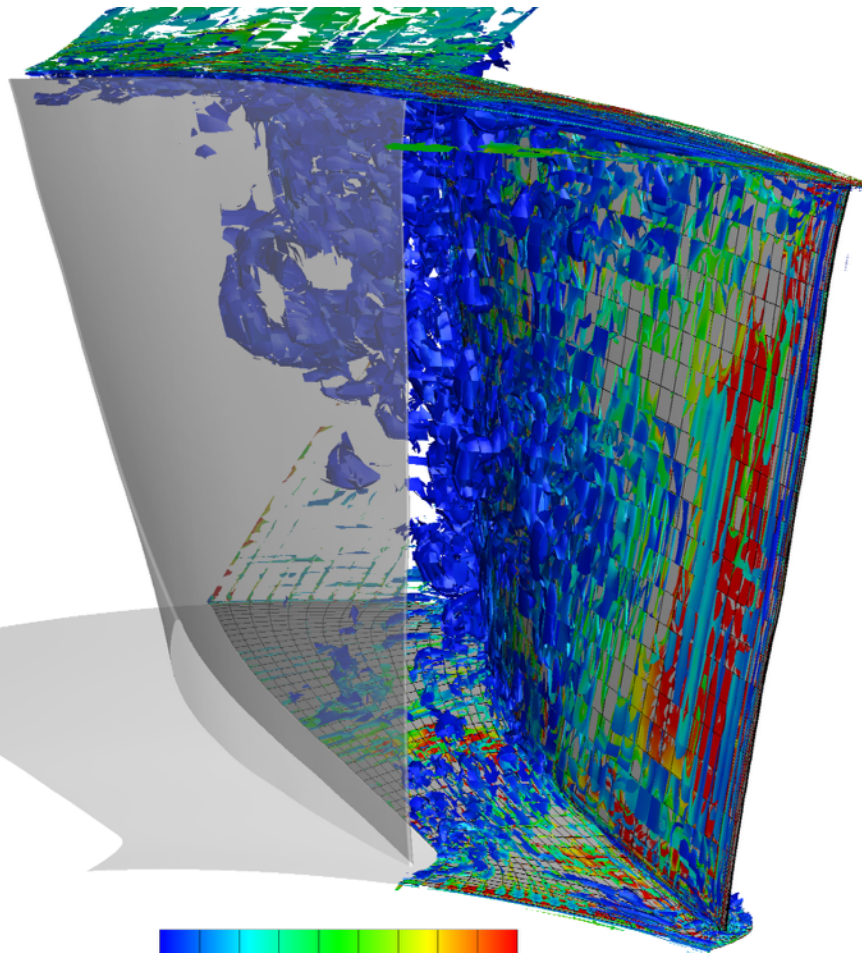


- Δ Exp.
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The profiles are in reasonable agreement with the experiments and the numerical results of Cahen et al.

X-LES of the transonic flow field in the NASA Rotor 37 (TC-P6 of TILDA)



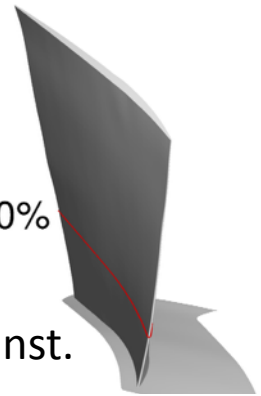
iso-surface of Q-criterion
160512 hexahedral elements

- P^2 computation using RO3-3
- Filter width $\Delta=5e-5$
- Boundary conditions
 - $p_{01} = 101325\text{Pa}$
 - $T_{01} = 288\text{K}$
 - $\omega = 1800\text{rad/s}$
 - $Tu_1 = 3\%$
 - $\alpha_1 = 0^\circ$
- parallel computation on 96 cores

According to our first experiences, for a practical usage of X-LES, **initializing with RANS** seems mandatory

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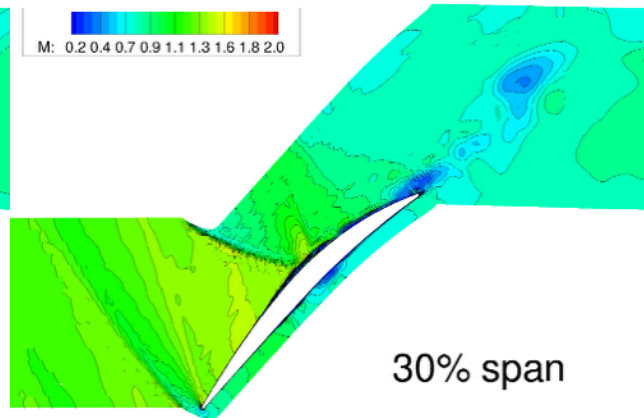
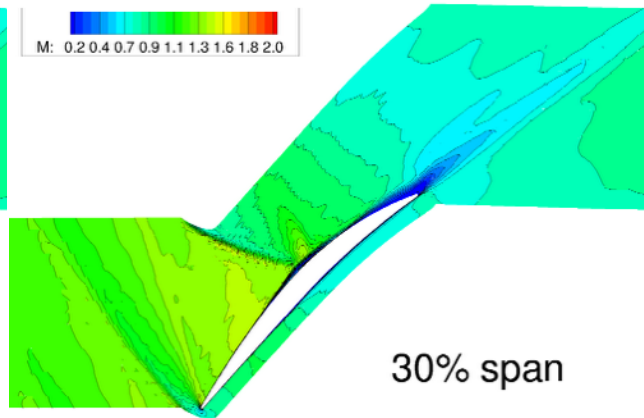
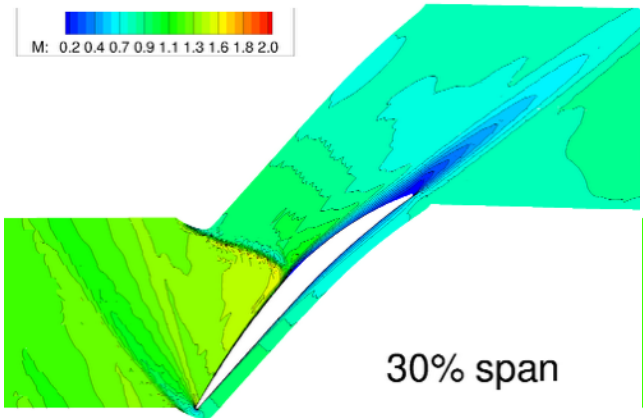
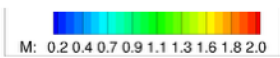
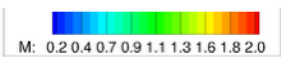
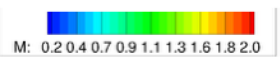
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RANS

X-LES ave.

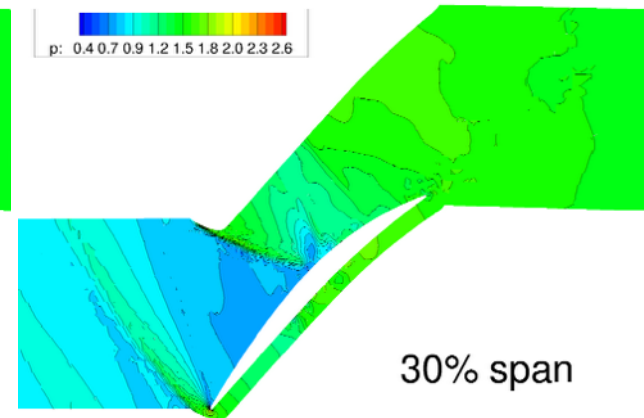
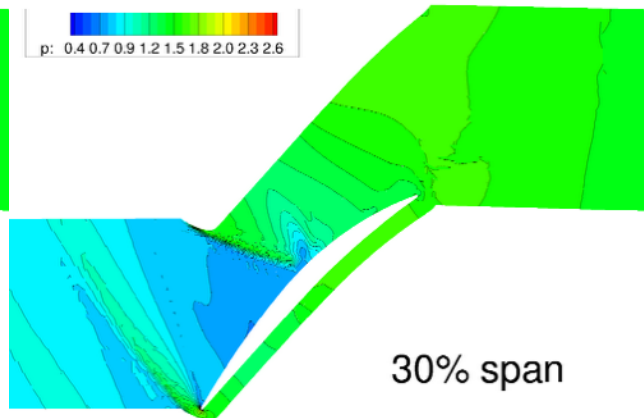
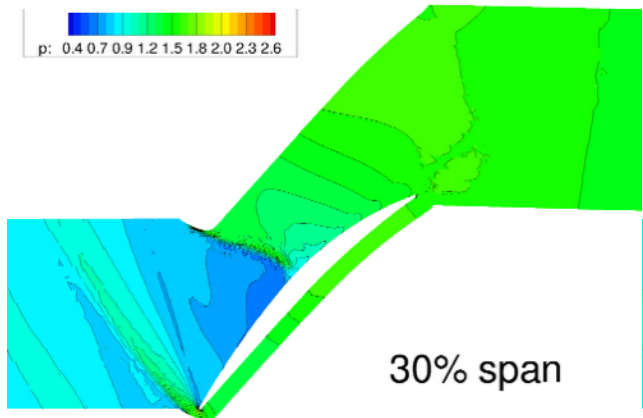
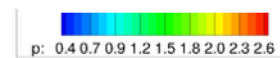
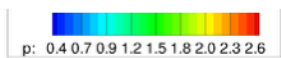
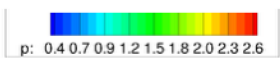
X-LES inst.



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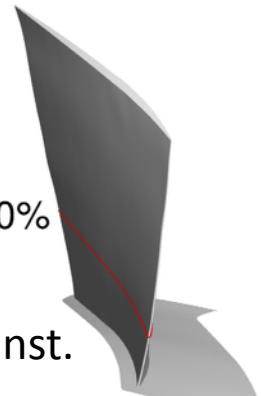
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Bassi et al. "Assessment of a high-order accurate Discontinuous Galerkin method for turbomachinery flows" *Accepted for publication on International Journal of CFD* (2016)

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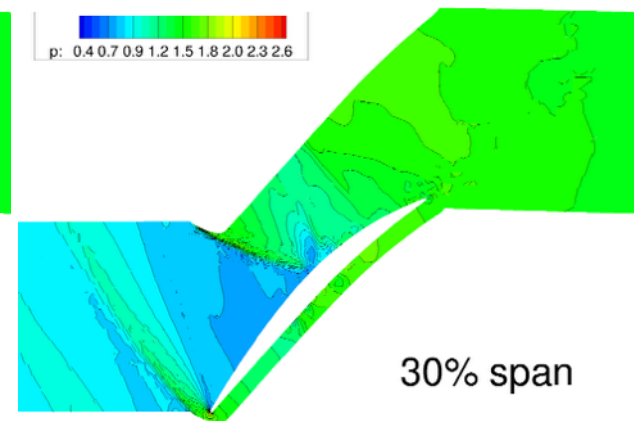
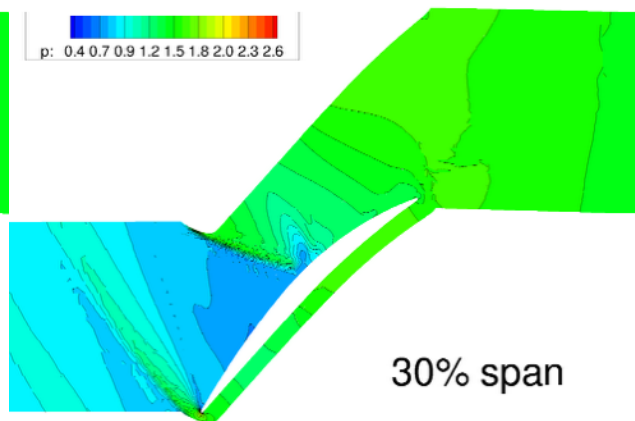
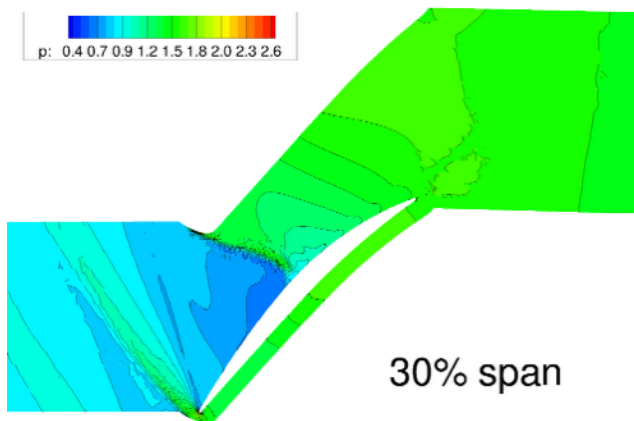
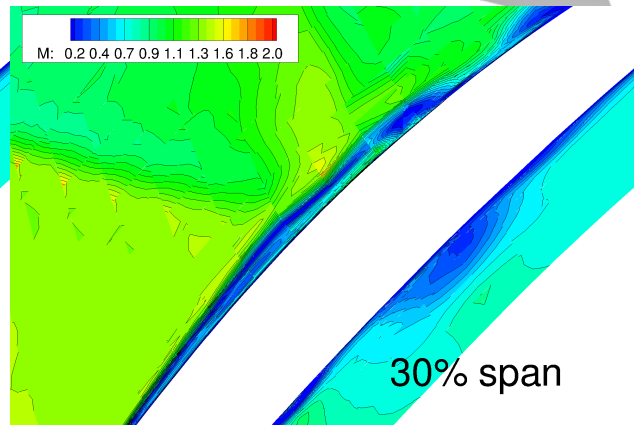
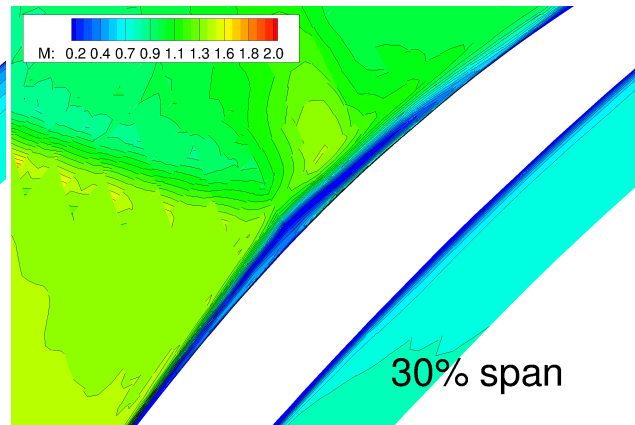
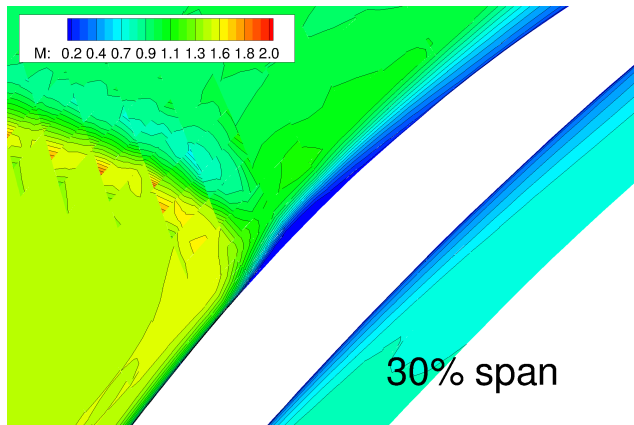
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RANS

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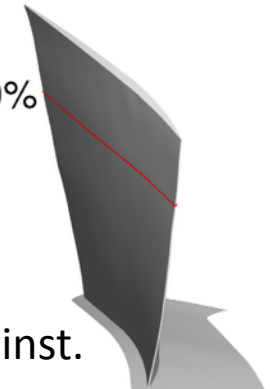
X-LES inst.



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X-LES of the transonic flow field in the NASA Rotor 37 (TC-P6 of TILDA)

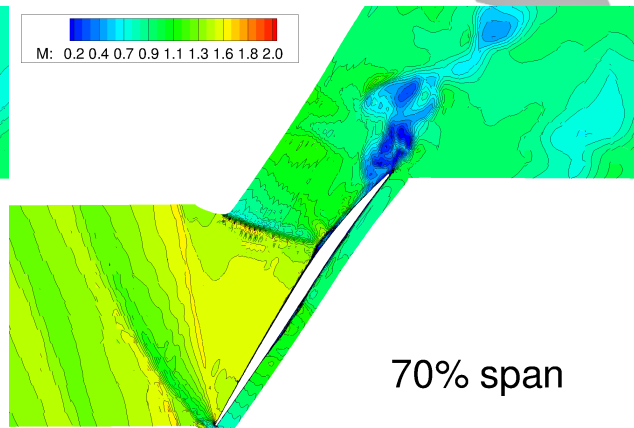
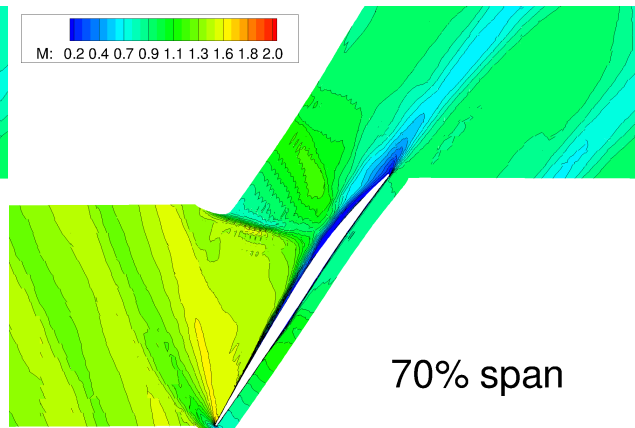
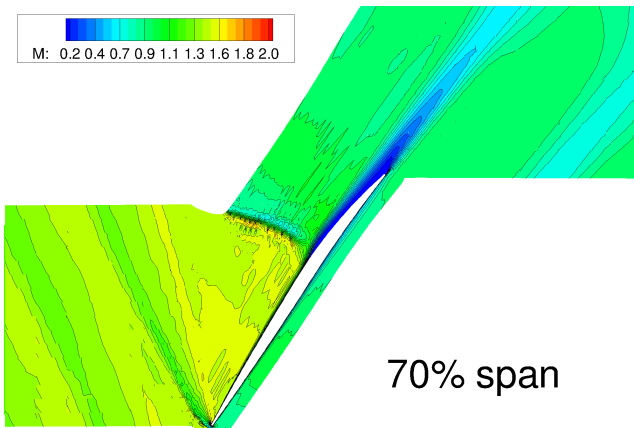
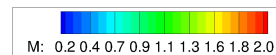
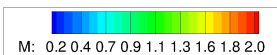
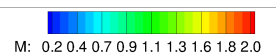
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RANS

X-LES ave.

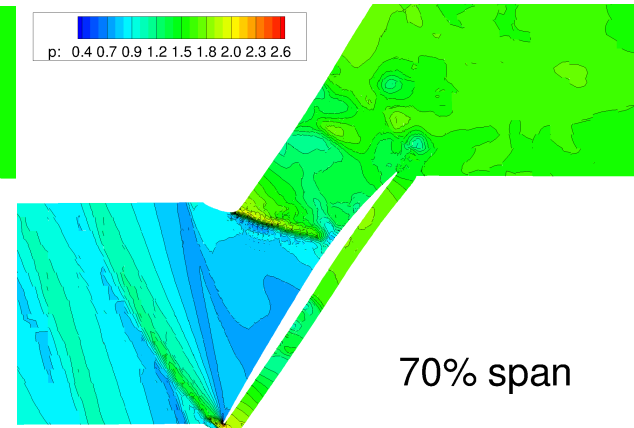
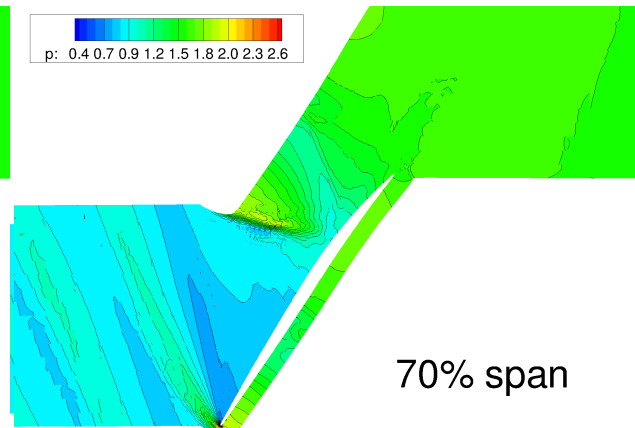
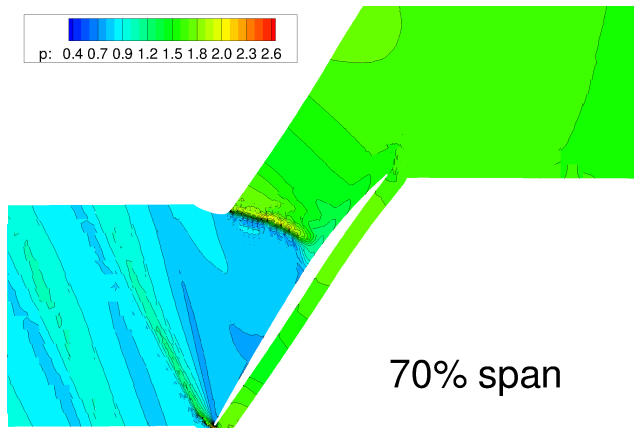
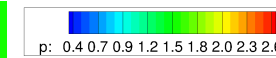
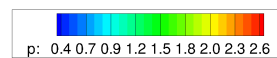
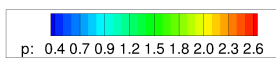
X-LES inst.



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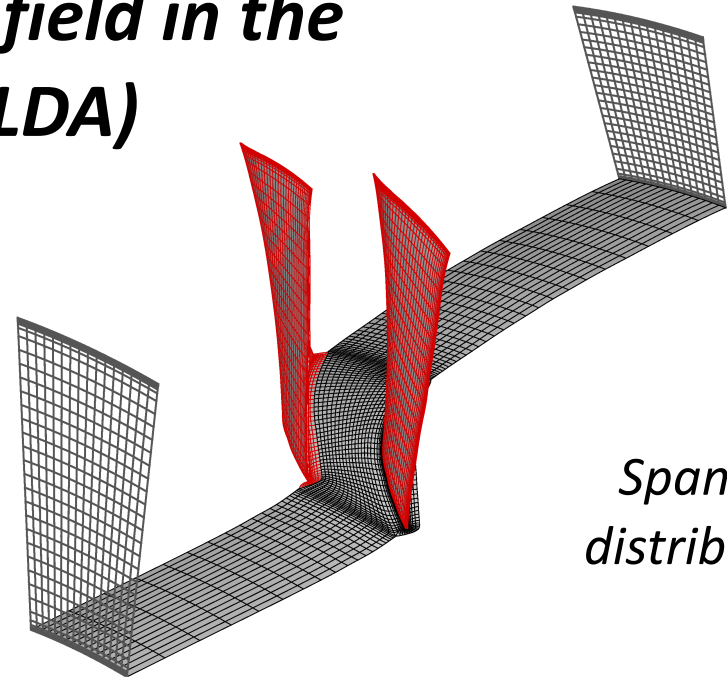
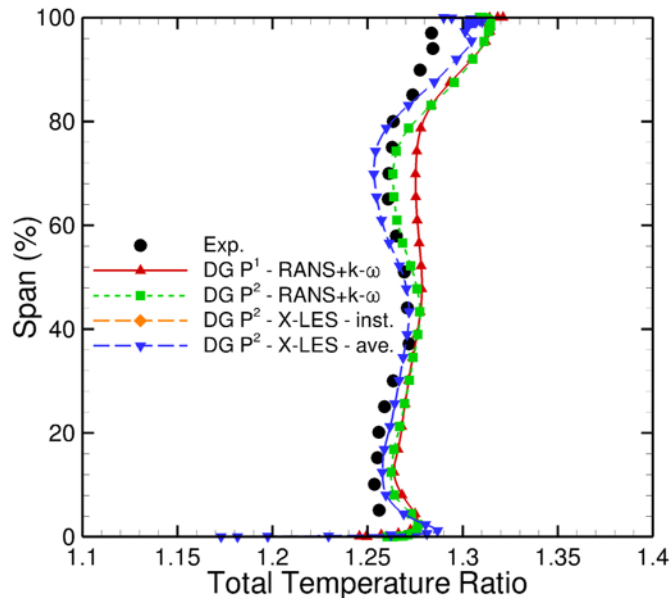
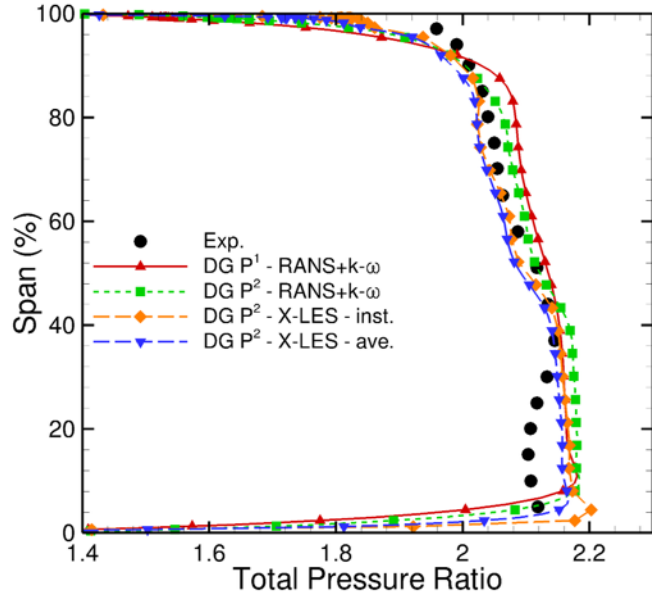
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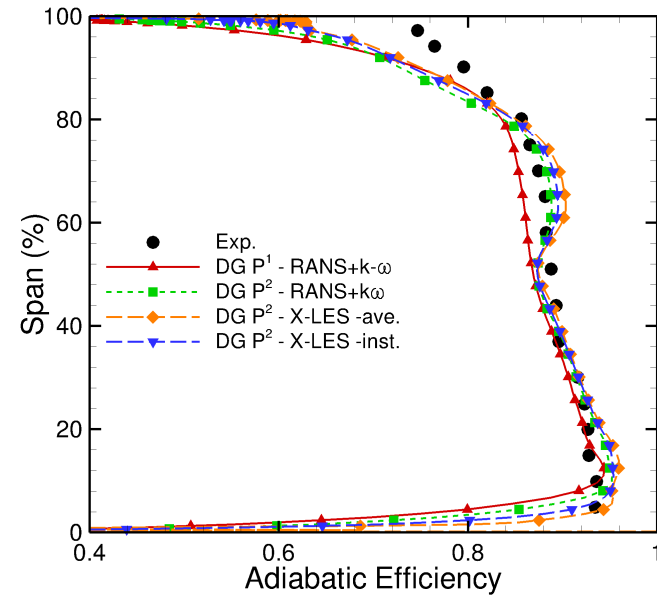
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Bassi et al. "Assessment of a high-order accurate Discontinuous Galerkin method for turbomachinery flows" *Accepted for publication on International Journal of CFD* (2016)

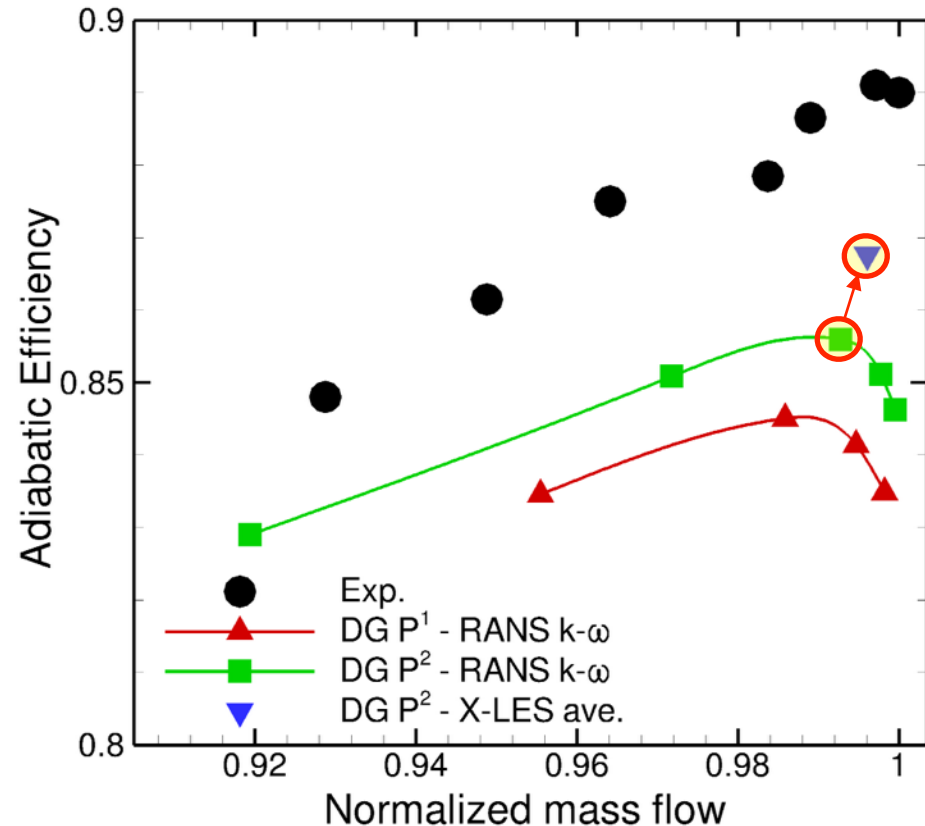
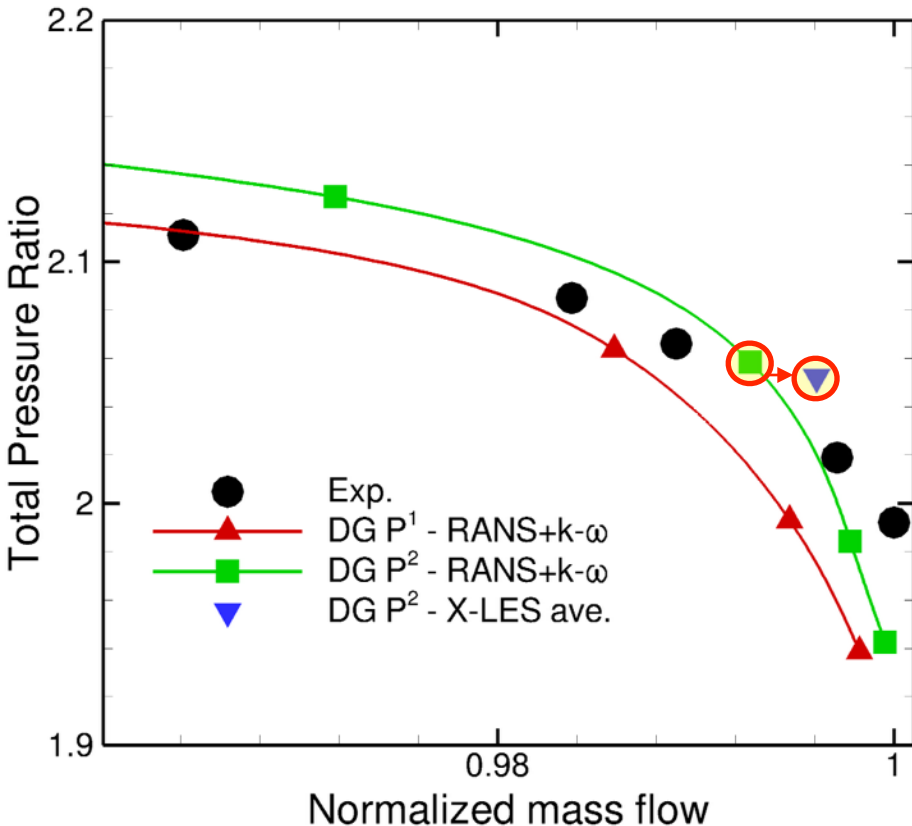
X-LES of the transonic flow field in the NASA Rotor 37 (TC-P6 of TILDA)



Spanwise distributions



X-LES of the transonic flow field in the NASA Rotor 37 (TC-P6 of TILDA)



Initializing the solution from a flow field corresponding to a normalized mass flow, resulting from a RANS computation, of ≈ 0.99
 X-LES moved towards ≈ 0.996

Conclusions...

- First results of an implicit DG solution of X-LES have been presented
- Implicit time accurate integration coupled with X-LES has shown to be robust

...and future work

- Compute the solution on a finer grid and/or a finer polynomial degree
- The influence on the solution of filter width parameter (Δ) need further investigation
- Evaluation of the numerical dissipation of the solver by computing the Decay of Isotropic Turbulence (DIT) test case

