Code MIGALE state-of-the-art

A. Colombo

HiOCFD4

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UNIVERSITÀ DEGLI STUDI DI BERGAMO

...with the contribution of

Francesco Bassi¹ Alessandro Colombo¹ Lorenzo Botti¹ Francesco Carlo Massa¹ Marco Savini¹ Nicoletta Franchina¹

Antonio Ghidoni² Gianmaria Noventa² Marco Lorini² Stefano Rebay²

Andrea Crivellini³

Carmine De Bartolo⁴ Alessandra Nigro⁴

Daniele Di Pietro⁵

Pietro Tesini⁶



¹ Università degli Studi di Bergamo
² Università degli Studi di Brescia
³ Università Politecnica delle Marche
⁴ Università della Calabria
⁵ University of Montpellier, France
⁶ SKF, Sweeden



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Brief code summary



- Discontinuous Galerkin (DG) method on hybrid grids
- Physical frame orthonormal basis functions
- 2D/3D steady and unsteady compressible and incompressible flows
- Explicit and implicit time accurate integration
- Fixed or rotating frame of reference
- Euler
- Navier–Stokes
- RANS coupled with the k- ω (EARSM)
- Hybrid RANS/LES (X-LES)
- MPI parallelism
- Fortran language

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Implicit accurate time integration

Several high-order temporal schemes are implemented

- Modified Extended BDF
- Two Implicit Advanced Step-point (TIAS)
- Explicit Singly Diagonally Implicit R-K (ESDIRK)
- linearly implicit Rosenbrock method





i) Hi-O schemes are more efficient than Lo-O ones for high required accuracy *ii)* Rosenbrock-type schemes are appealing both for accuracy and efficiency



Rosenbrock schemes in a nutshell (I/II)

From the DG spatial discretization we obtain a system of non-linear ODEs or DAEs

$$\mathbf{M}_{\mathbf{P}}(\mathbf{W}) \frac{d\mathbf{W}}{dt} + \mathbf{R}(\mathbf{W}) = \mathbf{0} \qquad \widetilde{\mathbf{R}} = \mathbf{M}_{\mathbf{P}}^{-1}\mathbf{R}$$
$$\mathbf{W}^{n+1} = \mathbf{W}^{n} + \sum_{j=1}^{s} m_{j}\mathbf{Y}_{j}$$
$$\left(\frac{\mathbf{M}_{\mathbf{P}}}{\gamma\Delta t} + \mathbf{J} - \frac{\partial\mathbf{M}_{\mathbf{P}}}{\partial\mathbf{W}}\widetilde{\mathbf{R}}\right)^{n}\mathbf{Y}_{i} = -\mathbf{M}_{\mathbf{P}}^{n}\left[\widetilde{\mathbf{R}}\left(\mathbf{W}^{n} + \sum_{j=1}^{i-1} a_{ij}\mathbf{Y}_{j}\right) - \sum_{j=1}^{i-1} \frac{c_{ij}}{\Delta t}\mathbf{Y}_{j}\right]$$
$$i = 1, \dots, s$$

only a linear system need to be solved for each stage

i.e. the Jacobian $J = \partial R / \partial W$ is assembled and factored only once per time step

With orthonormal basis functions (*physical space*) M_P reduces to the identity for compressible flows with conservative variables

For other sets of variables their DOFs can be coupled within $M_{\mathbf{P}}$ thus resulting in a matrix which can not be diagonal





Rosenbrock schemes in a nutshell (II/II)

Several Rosenbrock schemes, from order two to order six, have been compared

No need to "exactly" solve systems: GMRES tolerance can be increased with confidence with a significant reduction of WU

For a given order of accuracy, among the schemes considered, those with more stages are more accurate and efficient, *e.g.* RO5-8 vs. RO6-6

10⁻² 10⁻³ 10⁻³ 10⁻⁴ 10^{-4} 10⁻⁵ <u>10-6</u> <u>5</u> 10-7 ___10⁻⁵ ____10⁻⁶ 10⁻⁸ 10⁻⁷ RO2-2 10⁻⁹ RO3-3 10⁻⁸ RO3-4 **10**⁻¹⁰ RO4-6 10⁻⁹ **10**⁻¹¹ 305-8 306-6 R06-6 **10**⁻¹⁰ 10⁻¹² 3000 4000 5000 0.04 0.05 2000 0.03 0.06 0.07 0.08 1000 WU Λt

Bassi et al. "Linearly implicit Rosenbrock-type Runge-Kutta schemes for the Discontinuous Galerkin solution of compressible and incompressible unsteady flows" Computers & Fluids, 118 (2015) pp. 305 – 320





Working variables

Primitive variables to ensure the positivity of all thermodynamic variables at the discrete level....

We work with polynomial approximations not directly for p~ and ~T~ but for their logarithms $~\widetilde{p}=log(p)$ and $~\widetilde{T}=log(T)$

In this way the computed values $p = e^{\widetilde{p}}$ and $T = e^{\widetilde{T}}$ are always positive

Easy to implement: almost only change $\mathbf{M}_{\mathbf{P}}$...



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eXtra-Large Eddy Simulation (X-LES) in a nutshell (I/II)



Pros

- hybrid RANS\LES formulation independent from the wall distance
- use in LES mode of a clearly defined SGS based on the k-equation
- use of a k- ω turbulence model integrated to the wall

Cons

the *filter width* parameter is often related to the local element size

$$\begin{aligned} \frac{\partial}{\partial t}(\rho k) &+ \frac{\partial}{\partial x_j}(\rho u_j k) = \frac{\partial}{\partial x_j} \left[(\mu + \sigma^* \overline{\mu}_t) \frac{\partial k}{\partial x_j} \right] + P_k - D_k \\ \frac{\partial}{\partial t}(\rho \widetilde{\omega}) &+ \frac{\partial}{\partial x_j}(\rho u_j \widetilde{\omega}) = \frac{\partial}{\partial x_j} \left[(\mu + \sigma \overline{\mu}_t) \frac{\partial \widetilde{\omega}}{\partial x_j} \right] + (\mu + \sigma \overline{\mu}_t) \frac{\partial \widetilde{\omega}}{\partial x_k} \frac{\partial \widetilde{\omega}}{\partial x_k} \\ &+ P_\omega - D_\omega + C_D \end{aligned}$$

...an "original" interpretation for the X-LES implementation...

Bassi et al. "Time Integration in the Discontinuous Galerkin Code MIGALE - Unsteady Problems" In IDIHOM: Industrialzation of High-Order Methods - A Top-Down Approach, Vol. 128 of Notes on Numerical Fluid Me- chanics and Multidisciplinary Design Springer International Publishing

eXtra-Large Eddy Simulation (X-LES) in a nutshell (II/II)



$$\frac{\partial}{\partial t}(\rho k) + \frac{\partial}{\partial x_j}(\rho u_j k) = \frac{\partial}{\partial x_j} \left[(\mu + \sigma^* \overline{\mu}_t) \frac{\partial k}{\partial x_j} \right] + P_k - D_k$$
$$\overline{\mu}_t = \alpha^* \frac{\rho \overline{k}}{\hat{\omega}} \qquad D_k = \beta^* \rho \overline{k} \hat{\omega} \qquad \overline{k} = \max(0, k)$$

$$\hat{\omega} = \max\left(e^{\widetilde{\omega}_r}, \frac{\sqrt{\overline{k}}}{C_1 \Delta}\right)$$

$$\begin{array}{cccc} & \text{RANS} & \text{LES} & \text{ILES} \\ \hline \overline{\mu}_t & \alpha^* \frac{\rho \overline{k}}{e^{\widetilde{\omega}_r}} & \alpha^* \rho \sqrt{\overline{k}} C_1 \Delta & 0 \\ \hline D_k & \beta^* \rho \overline{k} e^{\widetilde{\omega}_r} & \beta^* \rho \frac{\overline{k}^{\frac{3}{2}}}{C_1 \Delta} & 0 \end{array}$$

X-LES of a shock BL interaction on a swept bump (AR2)



 P^2 converged computations with RANS+k- ω (also in its low-Re version) and EARSM1 have been performed and used as initialization for X-LES



72960 hexahedral elements

- Inlet boundary conditions
 - p_{0i} = 92000Pa
 - T_{0i} = 300K
 - Re_H = 1.69 x 10⁶
- Outlet static pressure used to impose the shock position (model dependent!)
- LBE to quickly find the "right" pressure ratio
- RO3-3 for the time-accurate solution
- Filter width Δ = 5e-2 (strong influence!)
 - 96 cores of our in-house Intel cluster :-(





The profiles are in reasonable agreement with the experiments and the numerical results of Cahen et al.





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X-LES of the transonic flow field in the NASA Rotor 37





- P² computation using RO3-3
- Filter width Δ =5e-5
- Boundary conditions
 - p₀₁ = 101325Pa
 - T₀₁ = 288K
 - ω= 1800rad/s
 - Tu₁ = 3%
 - $\alpha_1 = 0^\circ$
- still 96 cores of our in-house Intel cluster :-((

According to our first experiences, for a practical usage of X-LES, initializing with RANS seems mandatory

X-LES of the transonic flow field in the NASA Rotor 37





X-LES of the transonic flow field in the NASA Rotor 37







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Thursday, June 9 8:30-10:30 - MS 905 - 2 (Room 20) <u>Alessandra Nigro</u>, Carmine De Bartolo, Andrea Crivellini, Francesco Bassi MATRIX-FREE MODIFIED EXTENDED BACKWARD DIFFERENTIATION FORMULAE APPLIED TO THE DISCONTINUOUS GALERKIN SOLUTION OF COMPRESSIBLE UNSTEADY VISCOUS FLOWS

Thursday, June 9 14:30-16:30 - MS 905 - 3 (Room 20) Francesco Carlo Massa, <u>Gianmaria Noventa</u>, Francesco Bassi, Alessandro Colombo, Antonio Ghidoni, Marco Lorini HIGH-ORDER LINEARLY IMPLICIT TWO-STEP PEER METHODS FOR THE DISCONTINUOUS GALERKIN SOLUTION OF THE INCOMPRESSIBLE RANS EQUATIONS

Thursday, June 9 17:00-19:00 - MS 905 - 4 (Room 20) <u>Antonio Ghidoni</u>, Marco Lorini, Gianmaria Noventa, Francesco Bassi, Alessandro Colombo DISCONTINUOUS GALERKIN SOLUTION OF THE REYNOLDS- AVERAGED NAVIER–STOKES AND KL-KT-LOG(W) TRANSITION MODEL EQUATIONS

Friday, June 10 9:00-11:00 - MS 910 - 2 (Room 15)

Francesco Bassi, Lorenzo Botti, Alessandro Colombo, Andrea Crivellini, Antonio Ghidoni, Marco Lorini, Francesco Carlo Massa, Gianmaria Noventa ON THE IMPLEMENTATION OF X-LES IN A HIGH-ORDER IMPLICIT DG SOLVER

Tuesday, June 7 8:30-10:30 - CS 930 - 3 (Room 15)

Francesco Bassi, Alessandro Colombo, <u>Andrea Crivellini</u>, Matteo Franciolini HYBRID OPENMP/MPI PARALLELIZATION OF A HIGH–ORDER DISCONTINUOUS GALERKIN CFD SOLVER